

# Mode identification in rapidly rotating stars

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<sup>1</sup>University of Sheffield,

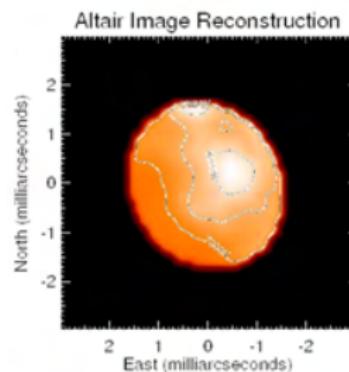
<sup>2</sup>High Altitude Observatory, Boulder

February 3, 2008



# Scientific context

- many intermediate and massive stars rotate rapidly
- their structure and evolution remain poorly understood



Monnier et al. (2007)

## CoRoT targets

| HD     | Type           | $v \cdot \sin i$   |
|--------|----------------|--------------------|
| 181555 | $\delta$ Scuti | 170 <sup>a</sup>   |
| 49434  | $\gamma$ Dor   | 90 <sup>a</sup>    |
| 171834 | $\gamma$ Dor   | 72 <sup>a</sup>    |
| 170782 | $\delta$ Scuti | 198 <sup>a</sup>   |
| 170699 | $\delta$ Scuti | > 200 <sup>a</sup> |
| 177206 | $\delta$ Scuti | > 200 <sup>b</sup> |

<sup>a</sup>Poretti et al. (2005), <sup>b</sup>Corotksy:  
[http://smsc.cnrs.fr/COROT/A\\_corotksy.htm](http://smsc.cnrs.fr/COROT/A_corotksy.htm).

## Recent 2D models

- Maeder & Meynet et al. (1997-2001)
- Roxburgh (2004, 2006)
- Self-Consistent Field (SCF) method: Jackson et al. (2004, 2005), MacGregor et al. (2007)
- ESTER project: Rieutord (2006), Espinosa & Rieutord (2007)

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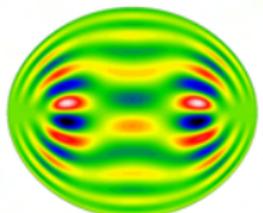
## Recent 2D calculations of acoustic modes

- Espinosa et al. (2004)
- Lovekin et al. (2008, 2009)
- Lignières et al. (2006), Reese et al. (2006, 2008, submitted to A&A)
- Lignières & Georgeot (2008, submitted to A&A)

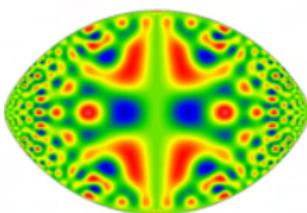
## Important results on the acoustic spectrum

- mode classification into regular and chaotic classes
  - polytropic case, using both ray dynamics and a normal mode analysis (Lignières & Georgeot, 2008, submitted to A&A)
  - extended to SCF models based on a normal mode analysis (Reese et al., submitted to A&A)

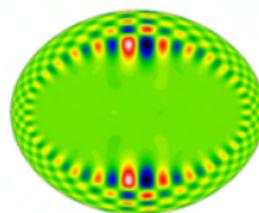
Island



Chaotic



Whispering gallery



$M=25.0M_{\odot}$     $\eta=0.6$     $\alpha=0.0$     $\omega=279.6\mu\text{Hz}$     $m=0^{\circ}$

$M=25.0M_{\odot}$     $\eta=0.9$     $\alpha=0.0$

$\omega=230.5\mu\text{Hz}$     $m=3^{\circ}$

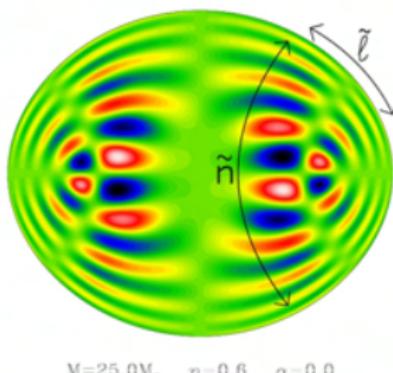
$M=25.0M_{\odot}$     $\eta=0.6$     $\alpha=0.0$

$\omega=313.6\mu\text{Hz}$     $m=1^{\circ}$

## Island modes

- rotating counterparts to modes with low  $\ell - |m|$  values
  - the most visible of the regular modes
  - new set of quantum numbers ( $\tilde{n}, \tilde{\ell}, m$ )
  - asymptotic formula valid for low azimuthal orders:

$$\omega = \tilde{n}\Delta_{\tilde{n}} + \tilde{\ell}\Delta_{\tilde{\ell}} + m^2\Delta_{\tilde{m}} - m\Omega + \tilde{\alpha}$$

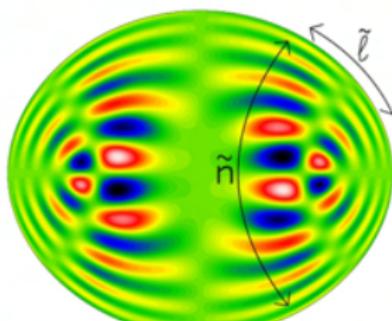


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- can this be used to identify pulsation modes?



# A new mode identification scheme

$$n_0 \leq \tilde{n} \leq n_1$$

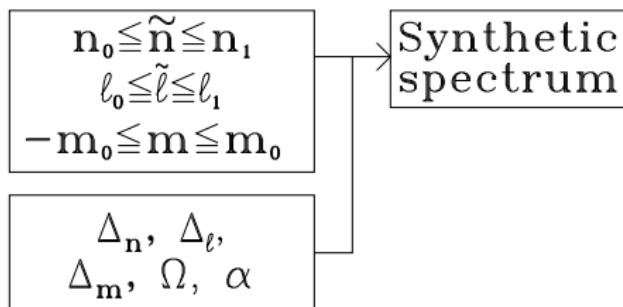
$$\ell_0 \leq \tilde{\ell} \leq \ell_1$$

$$-m_0 \leq m \leq m_0$$

$$\Delta_n, \Delta_\ell,$$

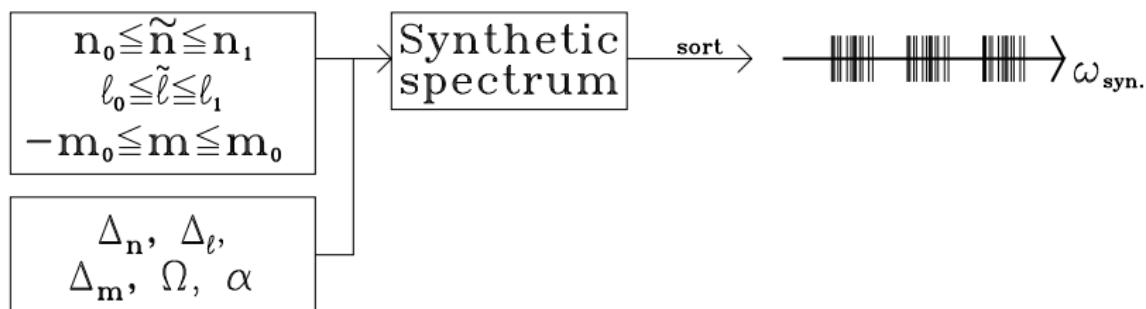
$$\Delta_m, \Omega, \alpha$$

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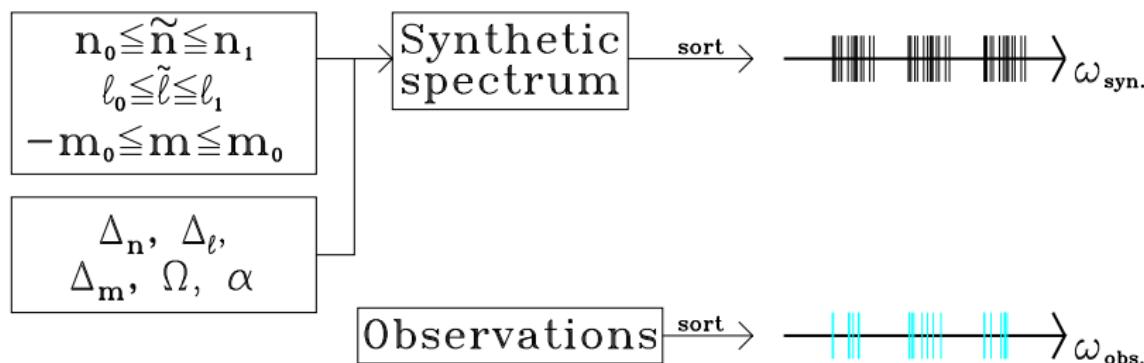
$$\omega = \tilde{n}\Delta_{\tilde{n}} + \tilde{l}\Delta_{\tilde{l}} + m^2\Delta_{\tilde{m}} - m\Omega + \tilde{\alpha}$$

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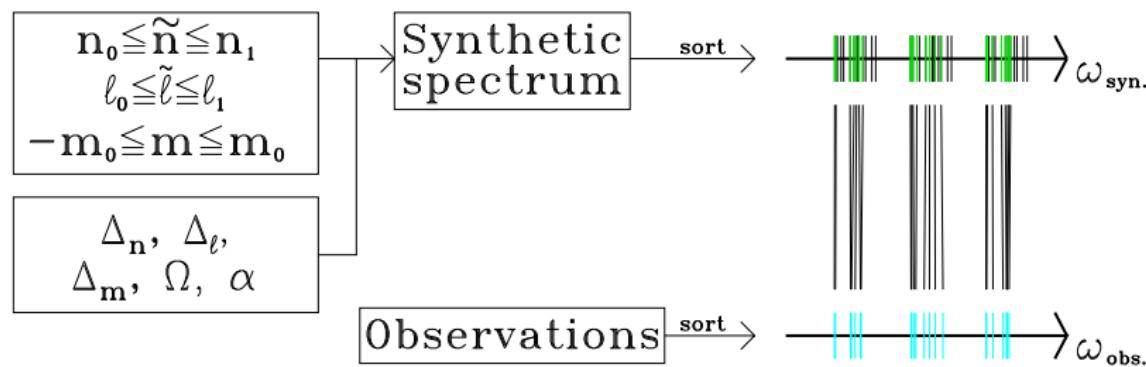
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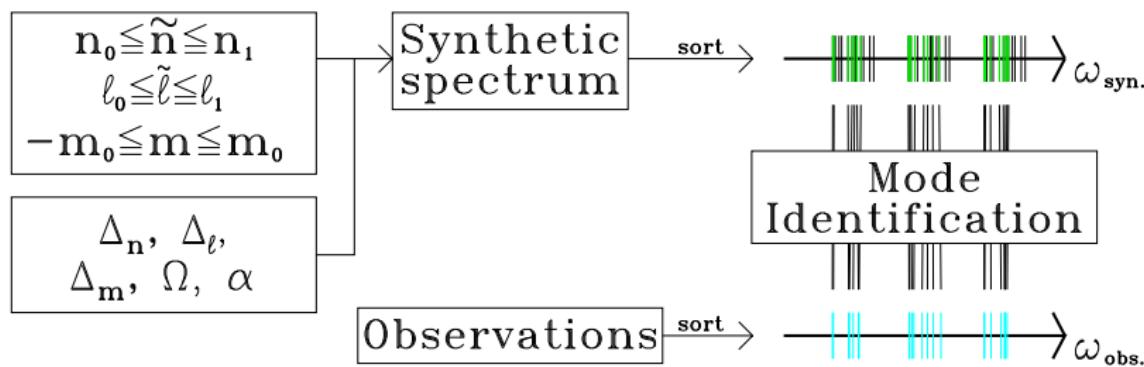
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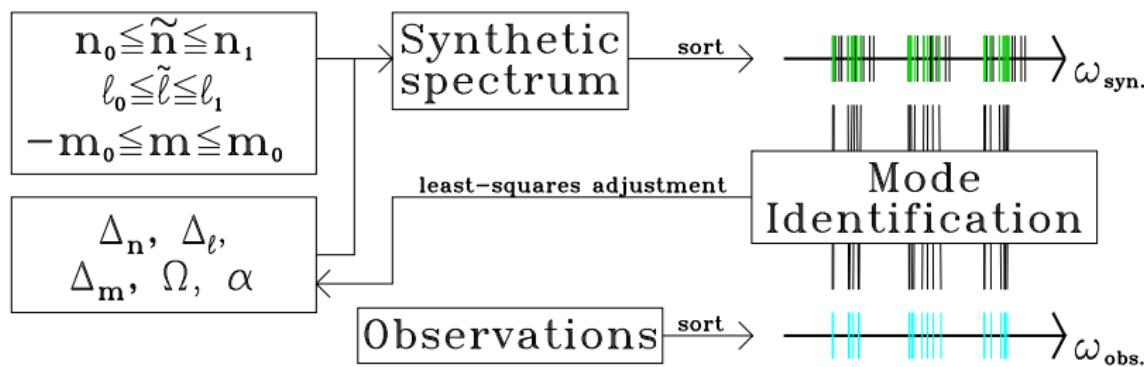
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- this yields a mode identification

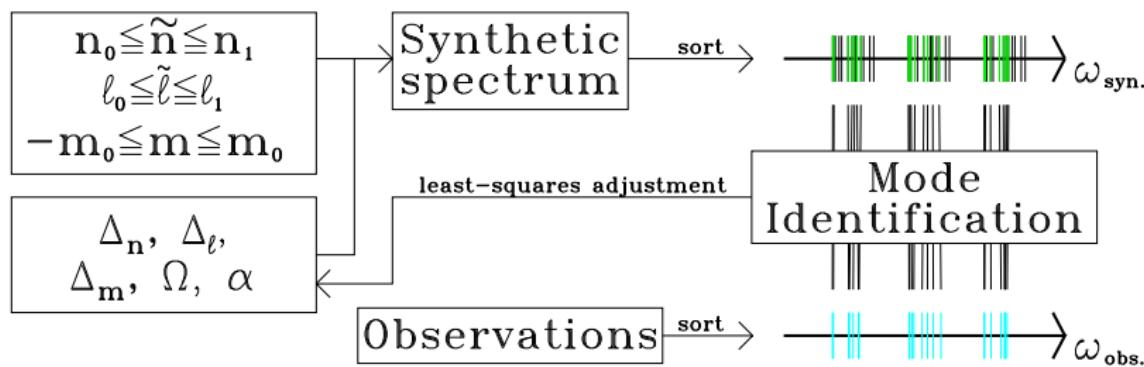
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- this yields a mode identification
- this can then be used to readjust  $\Delta_{\tilde{n}}$ ,  $\Delta_{\tilde{\ell}}$ ,  $\Delta_{\tilde{m}}$ ,  $\alpha$ ,  $\Omega$

# A new mode identification scheme



$$\omega = \tilde{n}\Delta_{\tilde{n}} + \tilde{\ell}\Delta_{\tilde{\ell}} + m^2\Delta_{\tilde{m}} - m\Omega + \tilde{\alpha}$$

- this yields a mode identification
- this can then be used to readjust  $\Delta_{\tilde{n}}, \Delta_{\tilde{\ell}}, \Delta_{\tilde{m}}, \alpha, \Omega$
- adjust parameters and obtain a closer fit to observed frequencies

## General procedure

- the previous slide explained how to find a plausible mode identification for a given set of parameters
- repeat the same procedure while scanning a 5D parameter space
- a slightly different set of parameters is used:

$$\Delta_{\tilde{n}}, \quad \frac{\Delta_{\tilde{\ell}}}{\Delta_{\tilde{n}}}, \quad \frac{\Delta_{\tilde{m}}}{\Delta_{\tilde{n}}}, \quad \frac{\alpha}{\Delta_{\tilde{n}}}, \quad \frac{\Omega}{\Delta_{\tilde{n}}}$$

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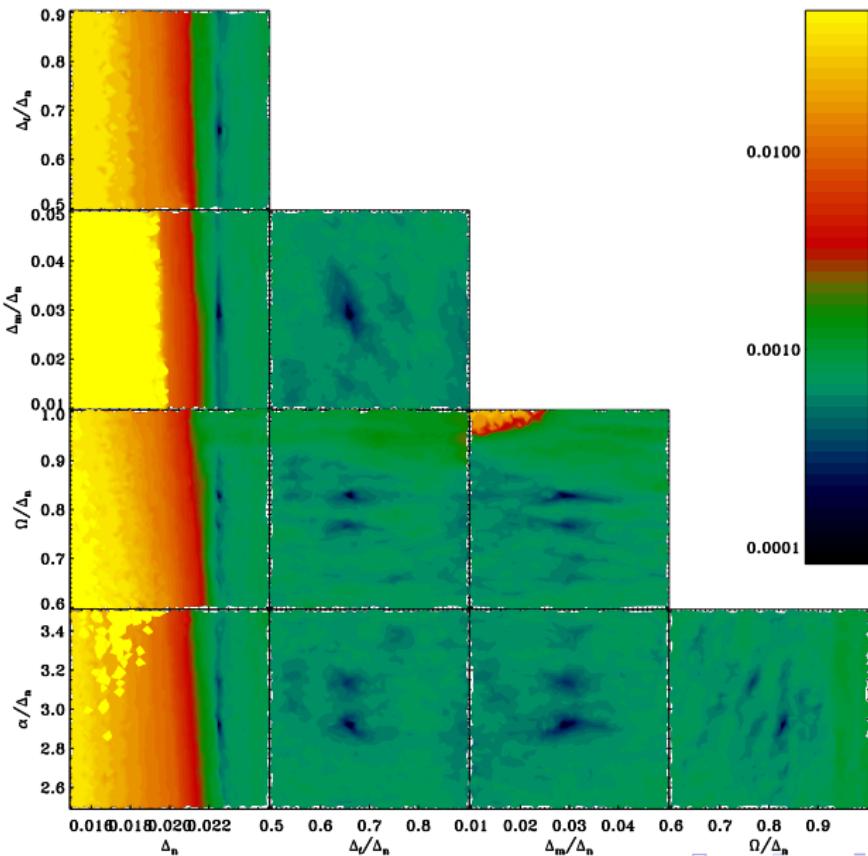
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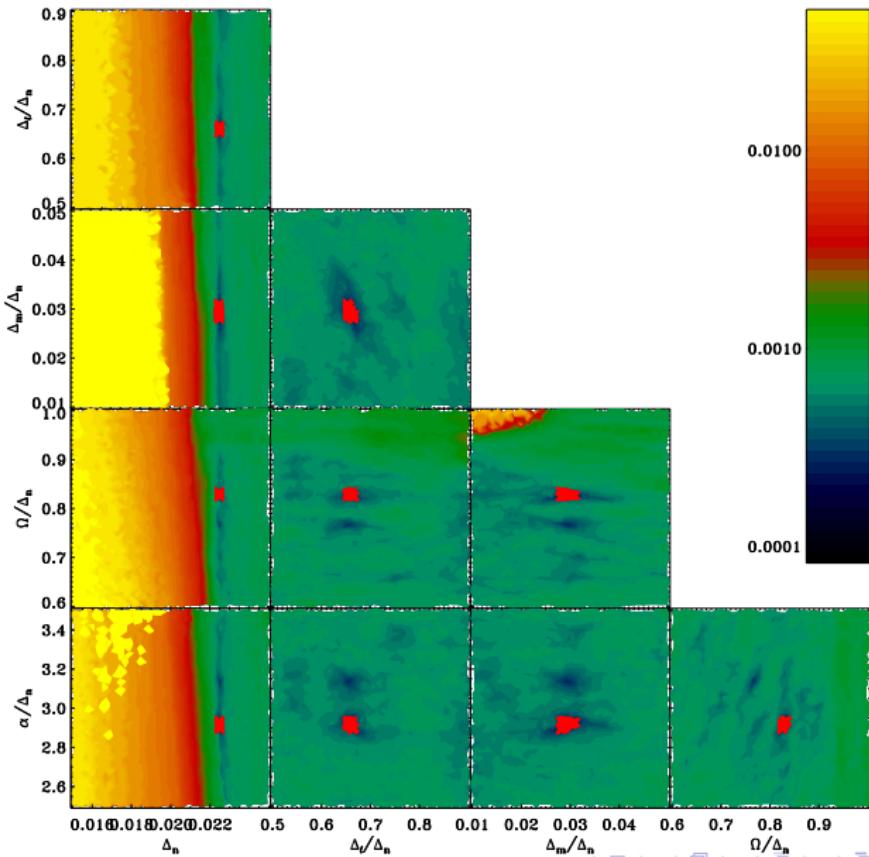
## Computational cost

- about 1 hour for  $50^5$  parameter combinations

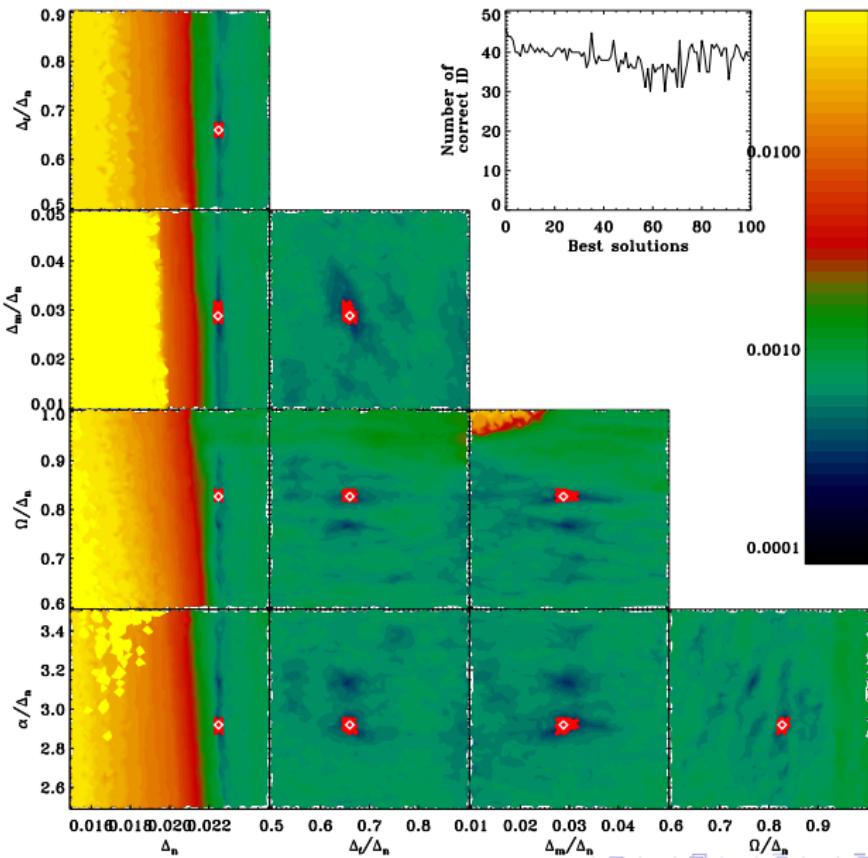
## Synthetic "Observations"



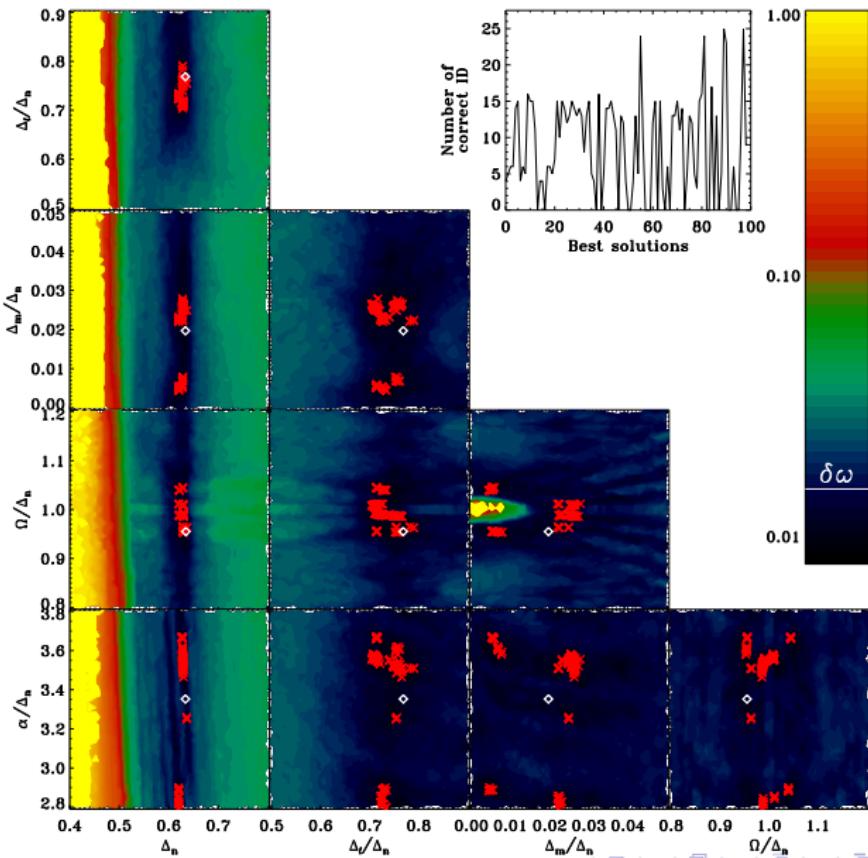
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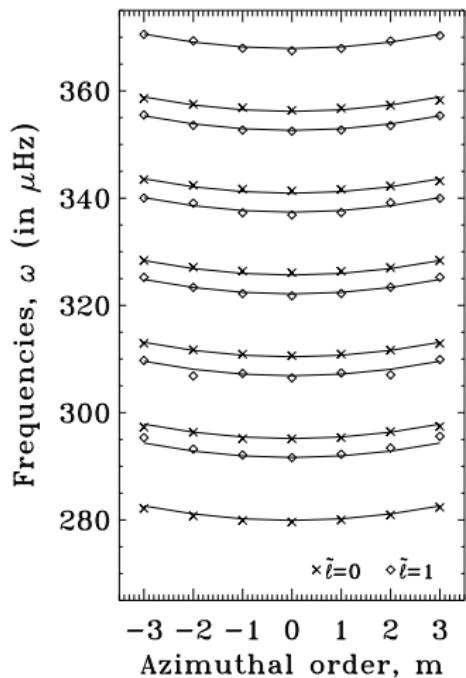


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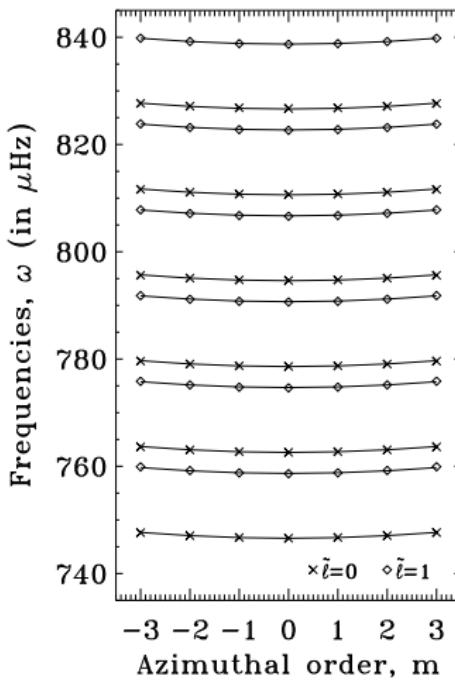
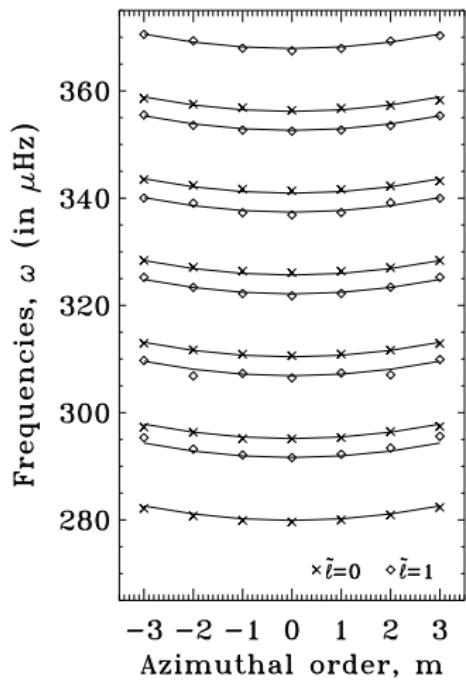


## Numerical frequencies



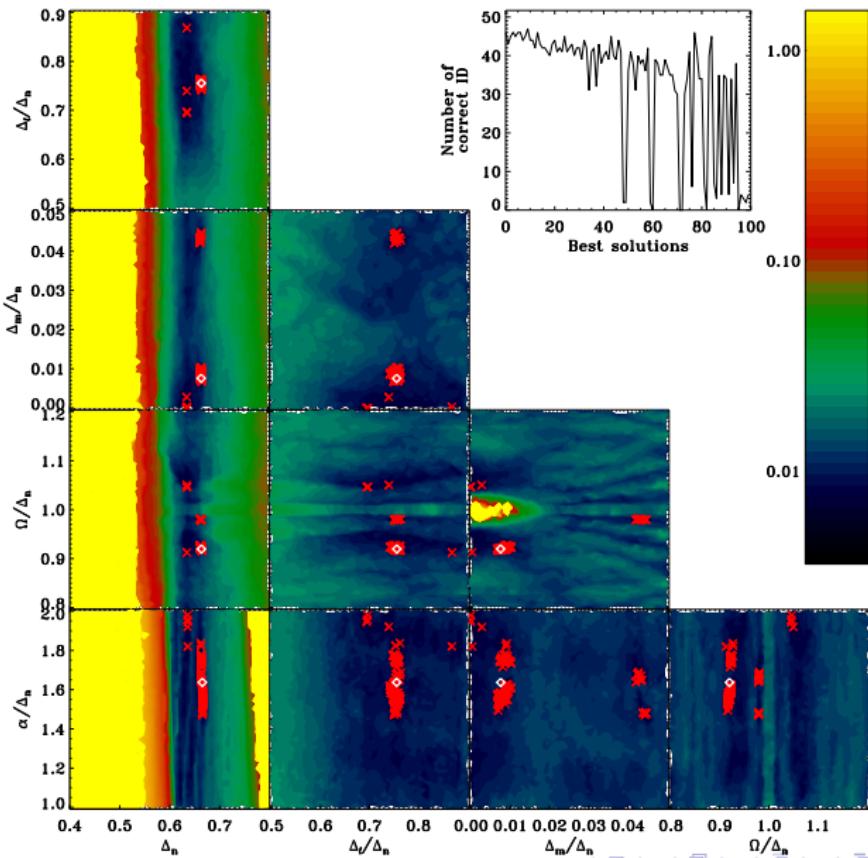


- avoided crossings cause deviations between asymptotic formula and numerical frequencies



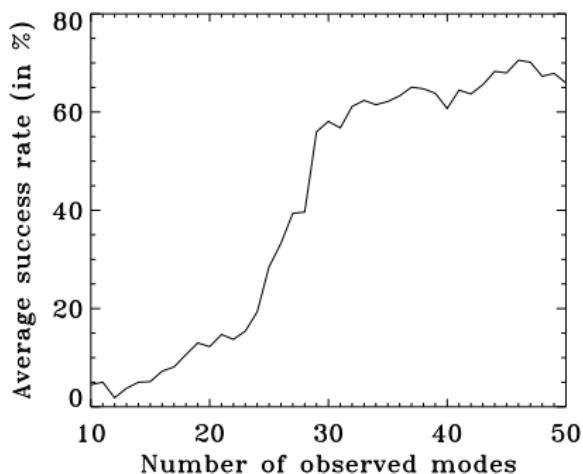
- avoided crossings cause deviations between asymptotic formula and numerical frequencies
- these deviations become much smaller at high radial orders

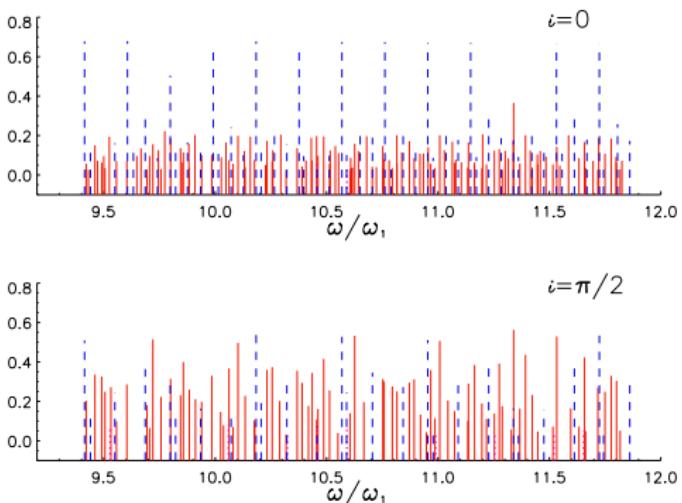
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- the previous example all contained 50 observed modes
- what happens when you change the number of modes?

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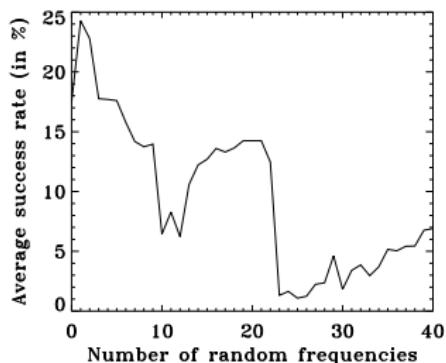


Lignières & Georgeot, 2009, submitted to A&A

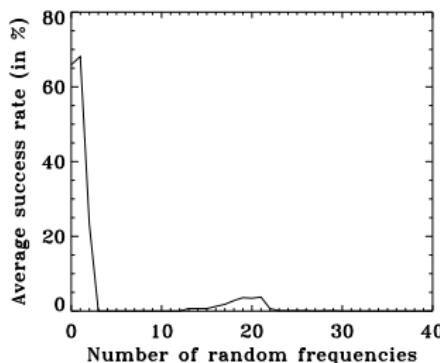
- chaotic modes are also likely to be visible in the frequency spectrum

- what happens to mode identification?

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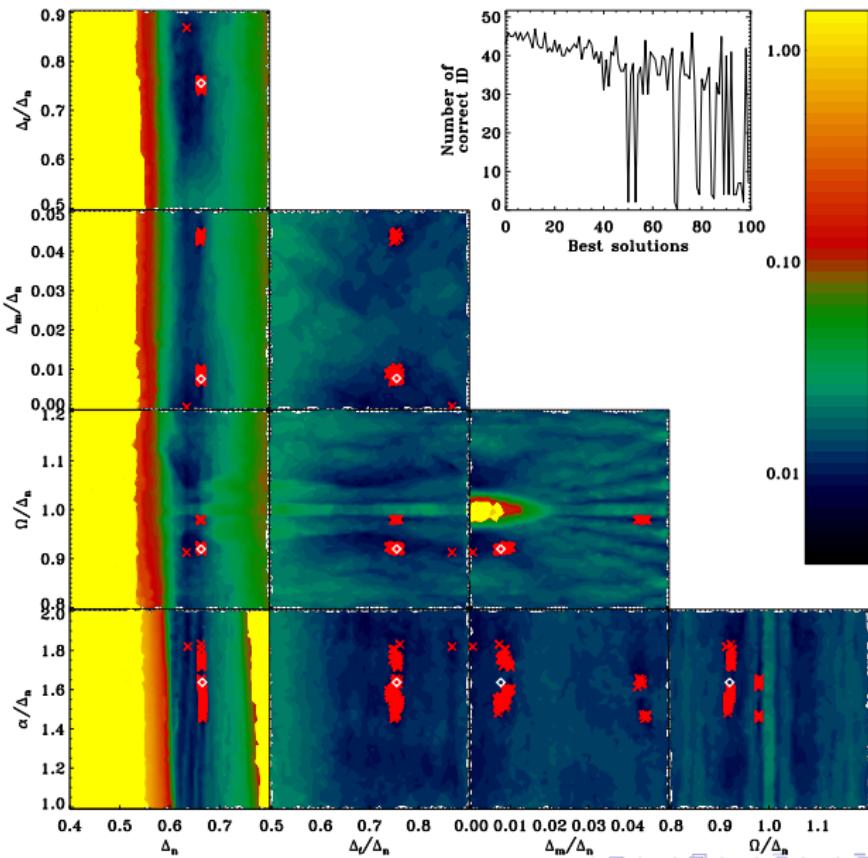


- low radial orders

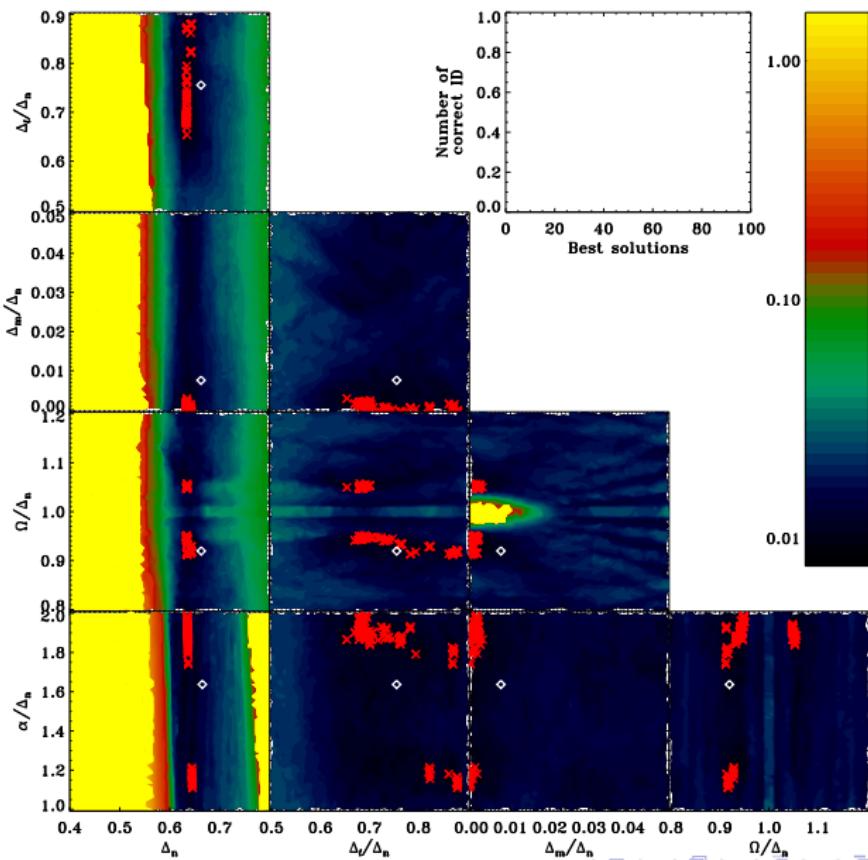


- high radial orders

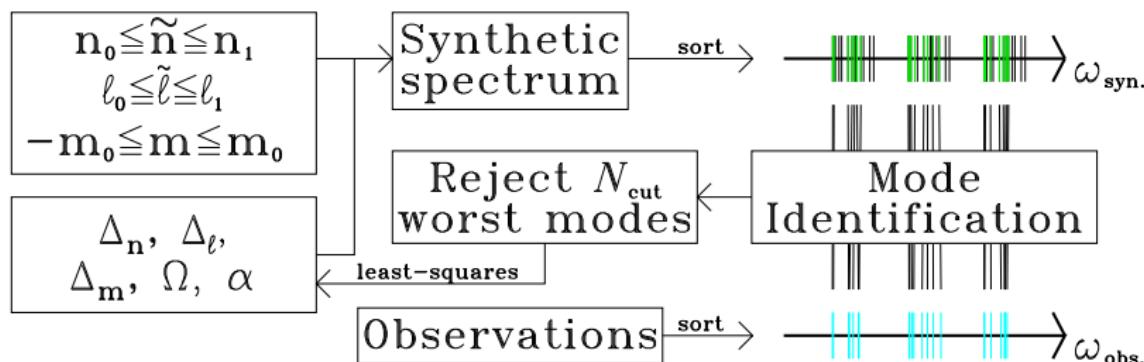
1 random frequency



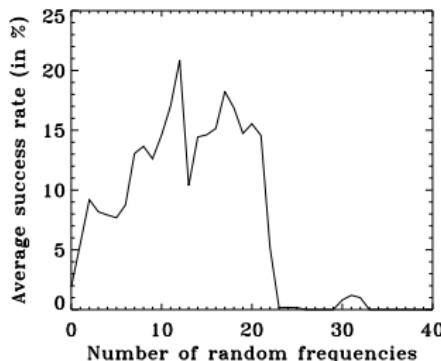
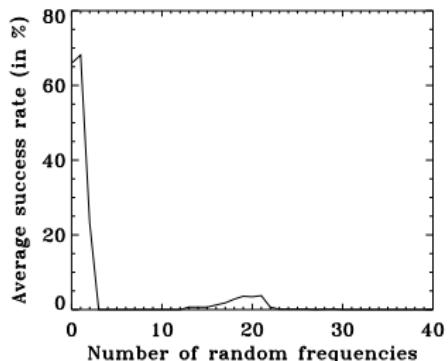
3 random frequencies



# Improving the mode identification scheme



- allow the algorithm to reject the  $N_{\text{cut}}$  worst solutions

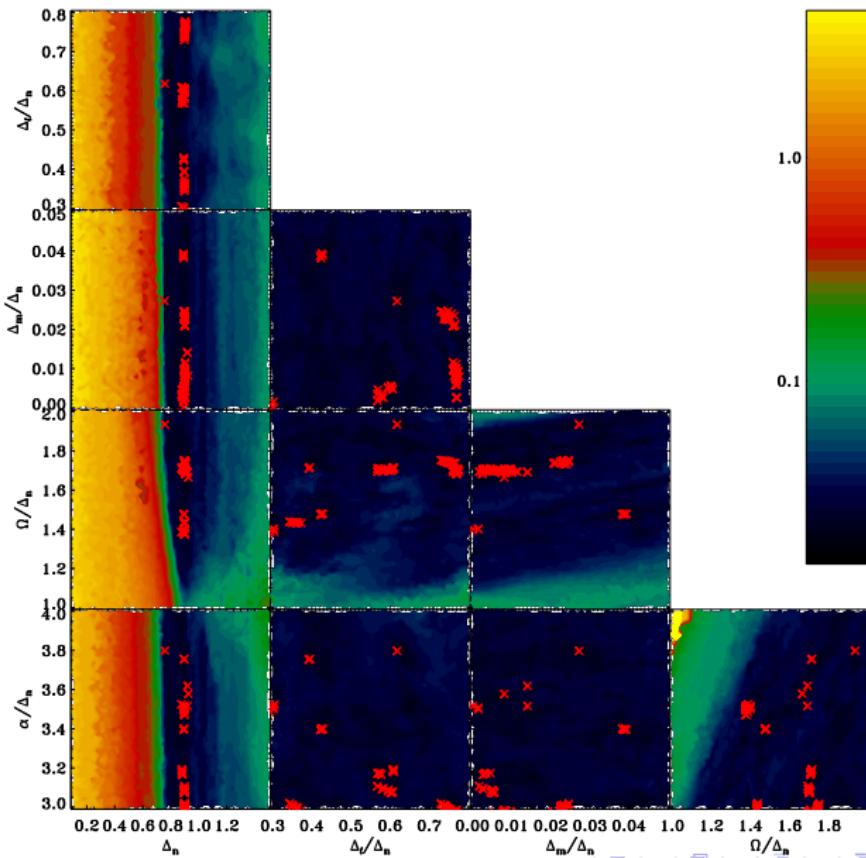


- keeping all frequencies
- rejecting 20 frequencies

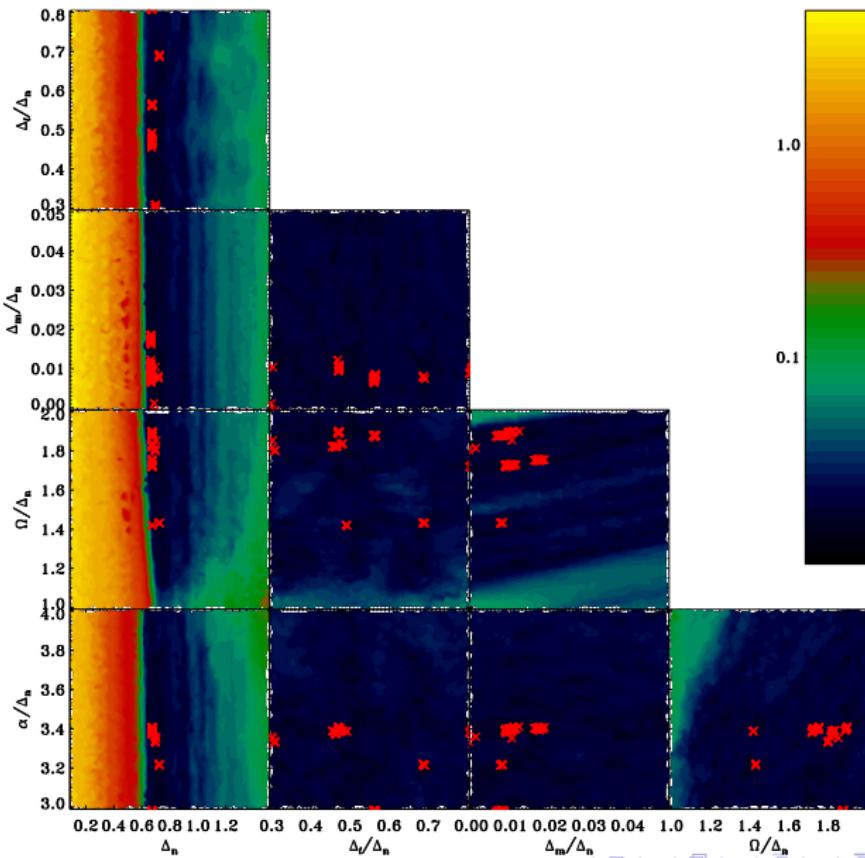
## Conclusion

- it is possible to identify acoustic modes in rapidly rotating stars based on the asymptotic formula for island modes
  - a sufficient number of modes is needed
  - modes with high radial orders yield better results
  - chaotic modes are problematic
- possible improvements:
  - correlate theoretical mode visibilities with observed mode amplitudes
  - use results from 2D pulsation calculations to restrict parameter space
  - take non-adiabatic effects into account to know which modes are excited
- direct comparison between observations and numerical calculations

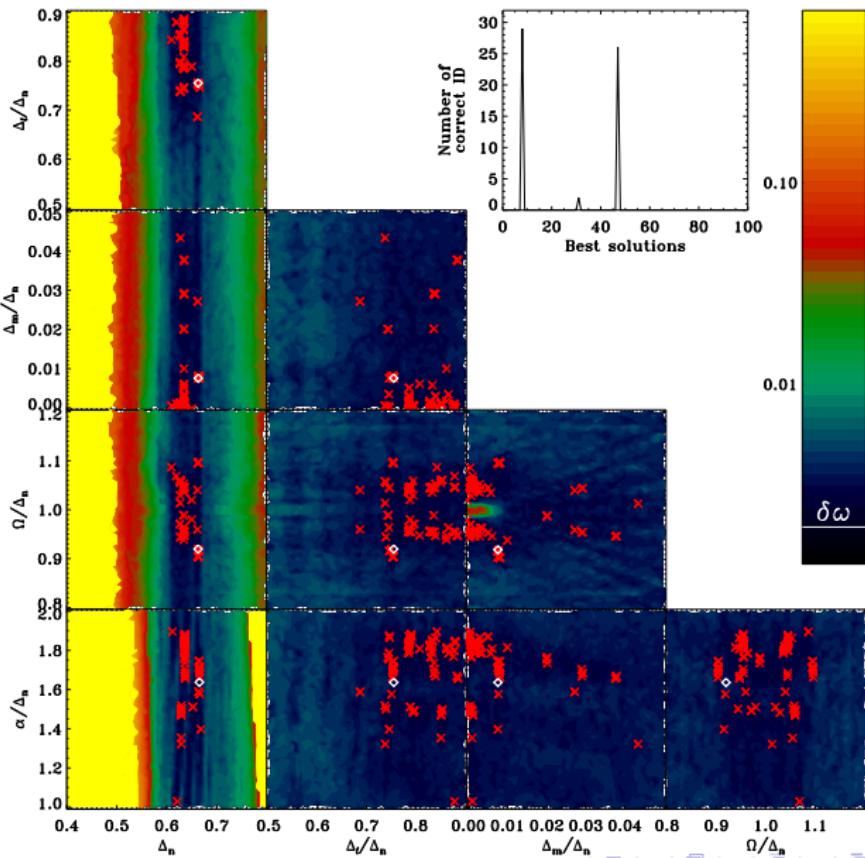
{ Ophiuchi: 19 frequencies



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$$N_{\text{cut}} = 20, N_{\text{rand}} = 0$$



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